

Deep Wavelet Scattering for Quantum Energy Regression

2016 March APS Meeting
Predicting and Classifying Materials via
High-Throughput Databases and Machine Learning I

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Collaborators

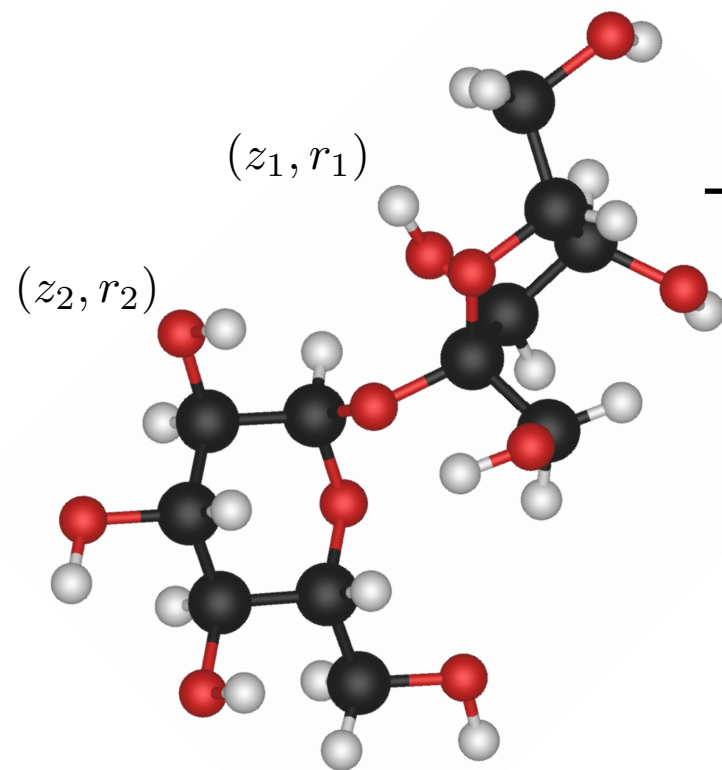


Stéphane Mallat



Nicolas Poilvert

What do we want do?



Molecule

compute →

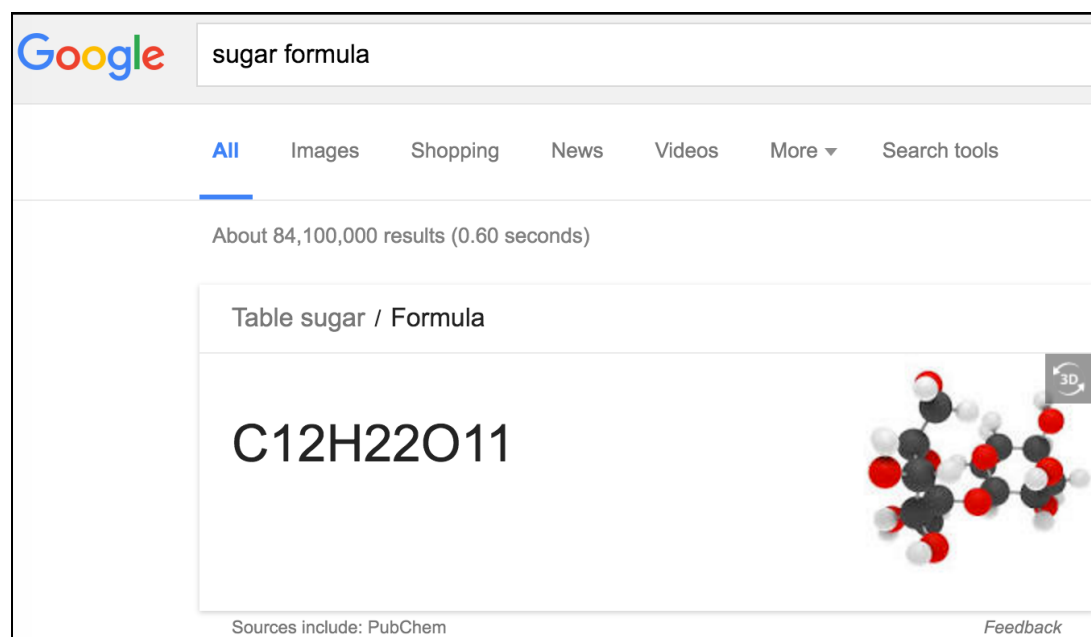
Ground state energy

$$f(x) = E_0(x) + \sum_{k>l} \frac{z_k z_l}{|r_k - r_l|}$$

Energy from
quantum electron
interactions

$$\begin{aligned} x &= \{\text{charges, positions}\} \\ &= \{(z_k, r_k)\}_k \end{aligned}$$

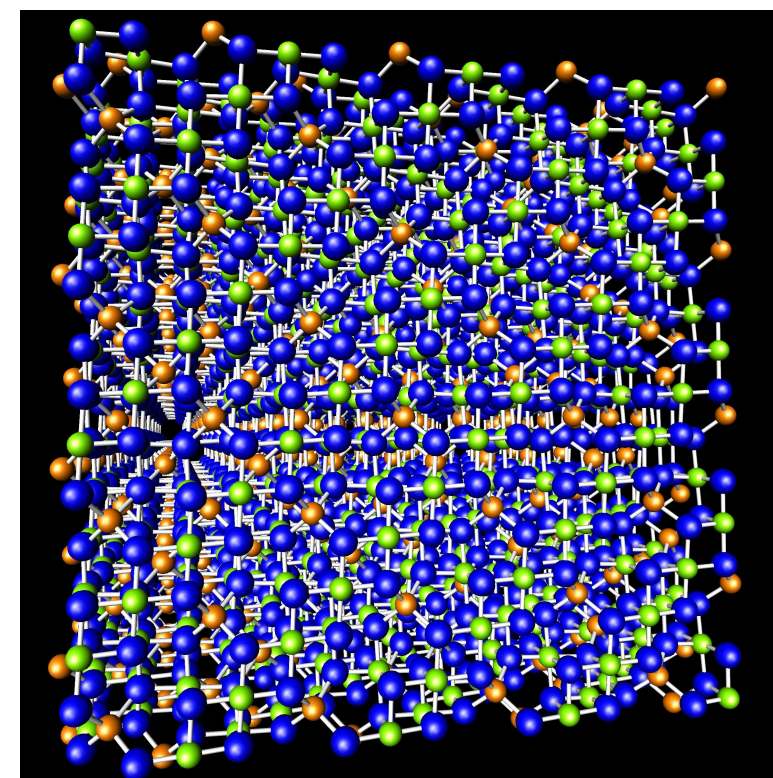
Why do we want to do it?



Google of molecules



Drug discovery



Materials design

So what's the problem?

$$f(x) = E_0(x) + \sum_{k>l} \frac{z_k z_l}{|r_k - r_l|}$$

Schrödinger equation

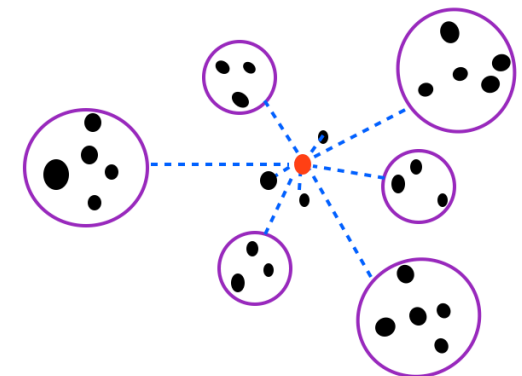
$$H[x]\Psi_0[x] = E_0(x)\Psi_0[x]$$

Not this

Fast multipole methods
(*Greengard, Rokhlin*)

Quantum mechanical (QM) approaches:

- Direct attacks (very small systems)
- Wave-function methods (small systems)
- Density functional theory (larger systems)



$O(N^\beta)$, $\beta \geq 3$
computational cost
 $N = \#$ electrons

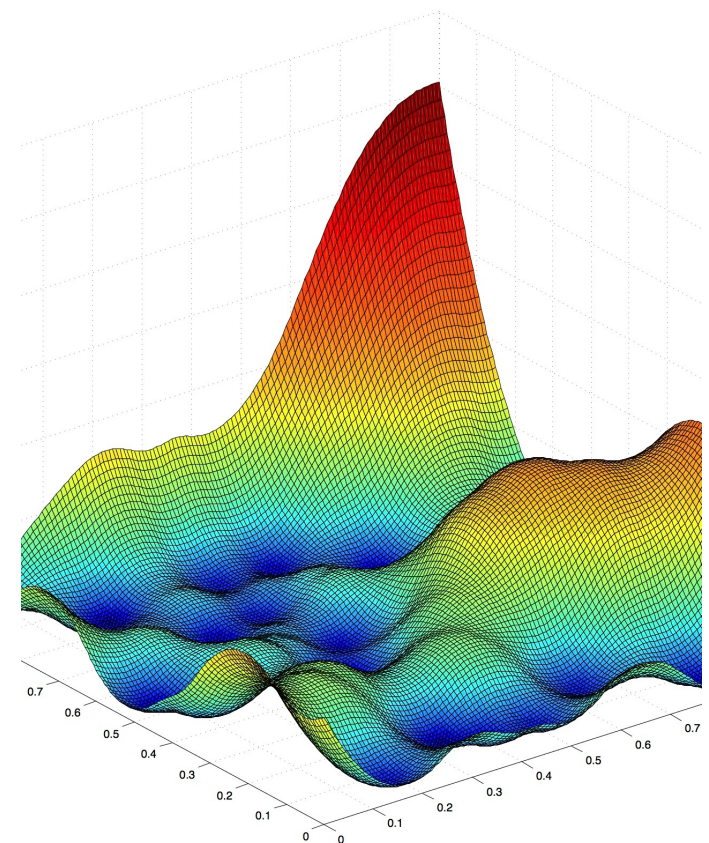
Recent idea:

Interpolation from known samples

- Use QM to compute training samples $\{(x_i, f(x_i))\}_{i \leq n}$
- Interpolate $f(x)$ from the training samples

Issue: Curse of dimensionality

- ϵ accuracy requires $O(\epsilon^{-d})$ samples
- $d = \text{dimension} = O(\# \text{ atoms})$



Recent idea:

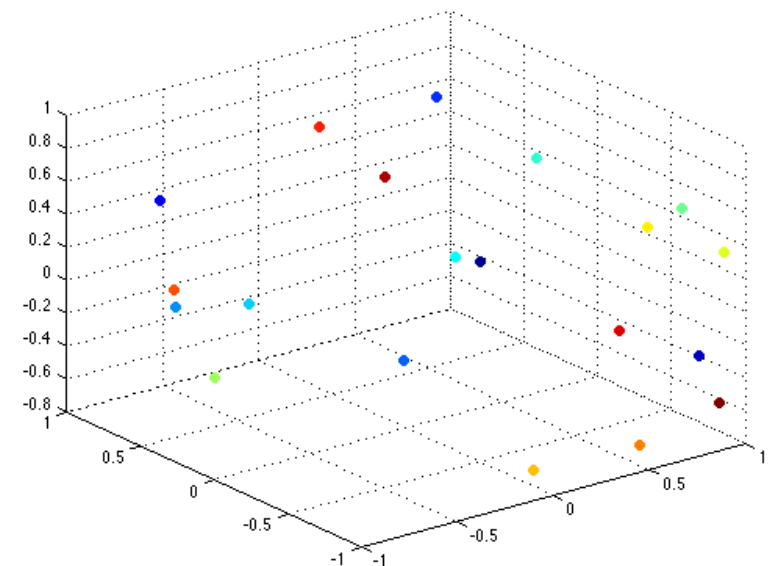
Interpolation from known samples

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Issue: Curse of dimensionality

- ϵ accuracy requires $O(\epsilon^{-d})$ samples
- $d = \text{dimension} = O(\# \text{ atoms})$

Need to take advantage of invariants and regularity of f



Energy invariants and regularity

- Permutation invariance

Invariant to permutations
of the atom index k

$$f(x) = E_0(x) + \sum_{k>l} \frac{z_k z_l}{|r_k - r_l|}$$

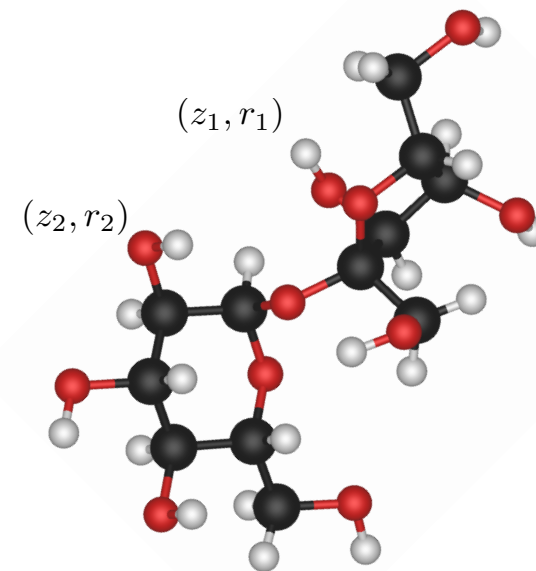
- Isometry invariance

Invariant to translations,
rotations and reflections

- Deformation stability

Lipschitz stable to
diffeomorphisms of
the molecule

(Bartók, Kondor, Csányi)



Molecule

$$\begin{aligned} x &= \{\text{charges, positions}\} \\ &= \{(z_k, r_k)\}_k \end{aligned}$$

Regression over a dictionary

- New representation of the molecule x :

$$\Phi(x) = (\phi_k(x))_k$$

- Linear regression (interpolation) of f using Φ :

$$\tilde{f}(x) = \langle w, \Phi(x) \rangle = \sum_k w_k \phi_k(x)$$

- Weights $(w_k)_k$ learned from the training data

(Bartók, Kondor, Csányi: SOAP kernels)

(Rupp, Tkatchenko, Müller, von Lilienfeld: Coulomb matrices)

(And many others...)

Dictionary properties

$$\tilde{f}(x) = \langle w, \Phi(x) \rangle = \sum_k w_k \phi_k(x)$$

The dictionary Φ should:

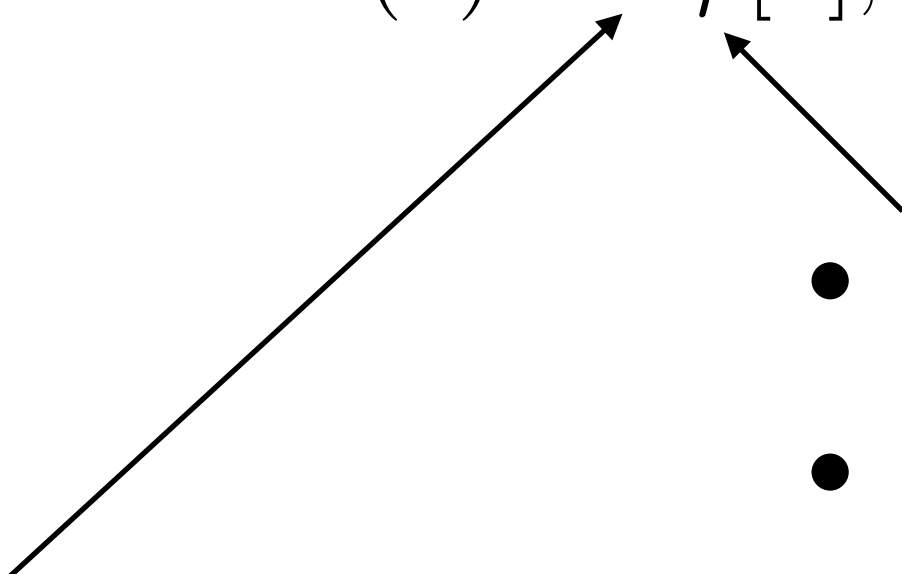
- Have the same invariants and regularity as f
 - permutation invariance
 - isometry invariance
 - deformation stability
- Span a large enough space to approximate f to high accuracy, with as few terms as possible

$$\|w\|_0 \leq M$$

Dictionary structure

Decompose Φ as:

$$\Phi(x) = \Theta \rho[x], \quad x = \{(z_k, r_k)\}_k$$

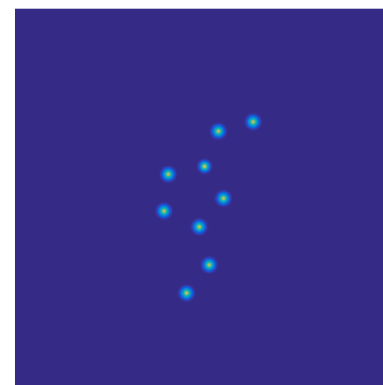
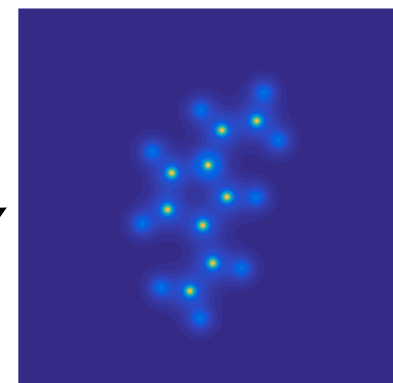
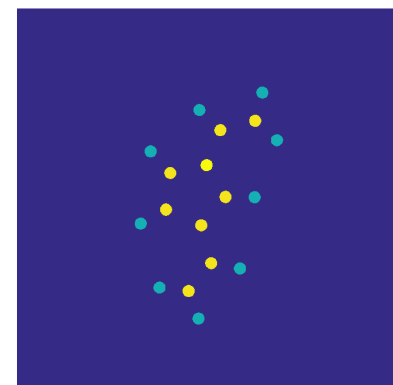
- 
- Isometry invariant
 - Lipschitz stable to diffeomorphisms
 - *This part is the main difficulty*
 - $\rho[x] : \mathbb{R}^3 \rightarrow \mathbb{R}$
 - Permutation invariant
 - Isometry covariant

Non-interacting density

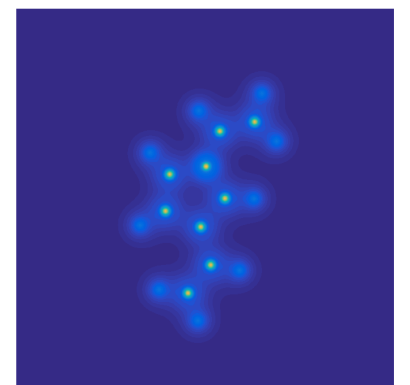
$$\rho[x](u) = \sum_k \rho[z_k](u - r_k), \quad x = \{(z_k, r_k)\}_k$$

Candidates for $\rho[z_k]$:

- $\rho[z_k] = z_k \delta$
- $\rho[z_k] = \text{Isolated atomic density}$
 $= \rho_{\text{cor}}[z_k] + \rho_{\text{val}}[z_k]$
- $\rho[z_k] = (\rho_{\text{cor}}[z_k], \rho_{\text{val}}[z_k])$



core



valence

Fourier and autocorrelation

Dictionary $\Phi = \Theta\rho$

Now we focus on Θ ; recall goals:

- Isometry invariant (translations, rotations, reflections)
- Stable to deformations

”Classic” translation invariant representations:

- Autocorrelation: $\Theta\rho(\tau) = \int \rho(u)\rho(u - \tau) du$
- Fourier modulus: $\hat{\Theta}\rho(\omega) = |\hat{\rho}(\omega)|^2$

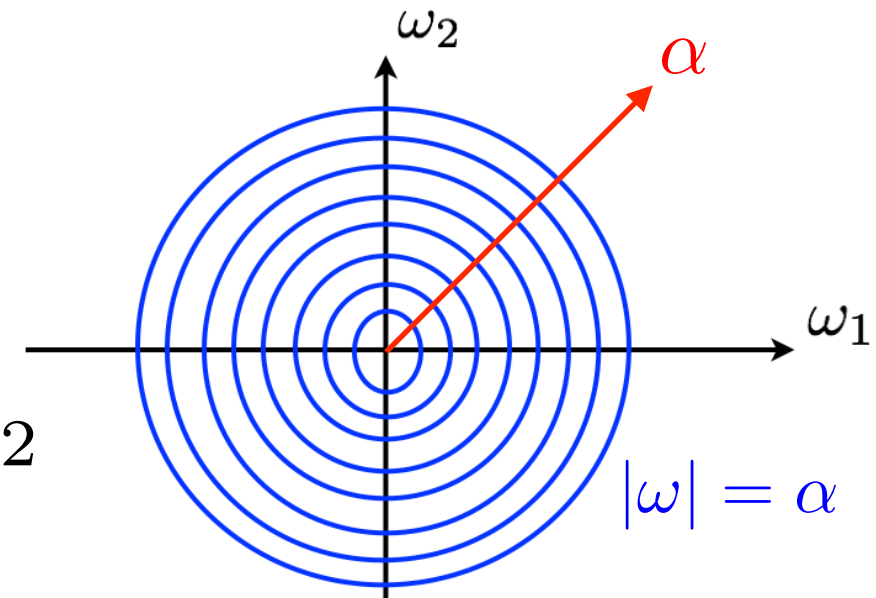
Integrate over rotations to obtain isometry invariance

Fourier and autocorrelation

Invariant Fourier operator:

$$\hat{\Theta}\rho(\alpha) = \|\hat{\rho}_\alpha\|_2^2 = \int_{S^2} |\hat{\rho}(\alpha\eta)|^2 d\eta$$

$$\omega = \alpha\eta, (\alpha, \eta) \in \mathbb{R}^+ \times S^2$$



Pros:

- Isometry invariant
- Diagonalizes Coulomb:

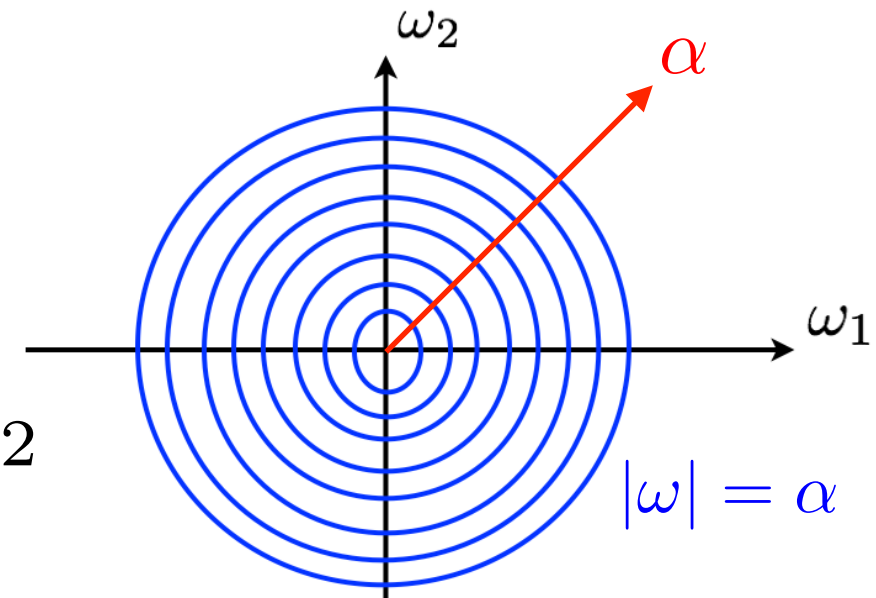
$$U(\rho) = \iint \frac{\rho(u)\rho(v)}{|u-v|} du dv = \frac{1}{2\pi^2} \int_0^\infty \alpha^{-2} \|\hat{\rho}_\alpha\|_2^2 d\alpha$$

Fourier and autocorrelation

Invariant Fourier operator:

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$$\omega = \alpha\eta, (\alpha, \eta) \in \mathbb{R}^+ \times S^2$$



Cons:

- Not sparse: $U(\rho) = \sum_{k=1}^{\epsilon^{-2}} w_k \|\hat{\rho}_{k\epsilon}\|_2^2 + O(\epsilon)$

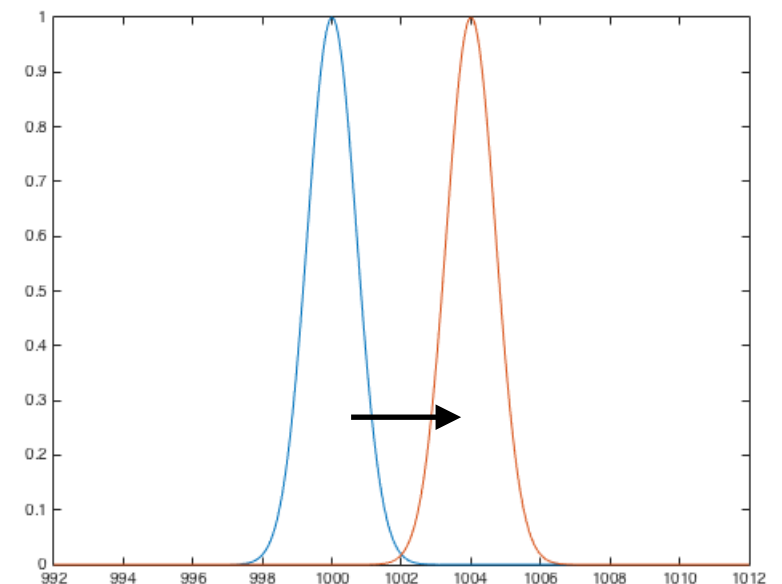
- Not stable to deformations:

Invariant autocorrelation:

– Bumps located at $\{|r_k - r_l|\}_{k,l}$

– Δ sized diffeomorphism

moves them $\Delta|r_k - r_l|$ distance

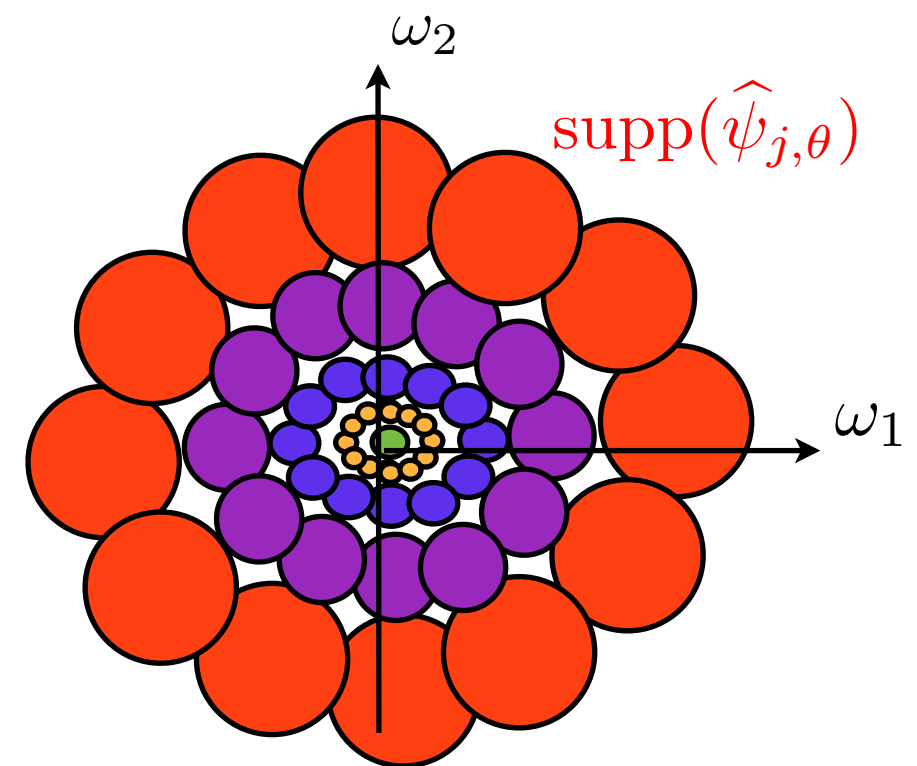
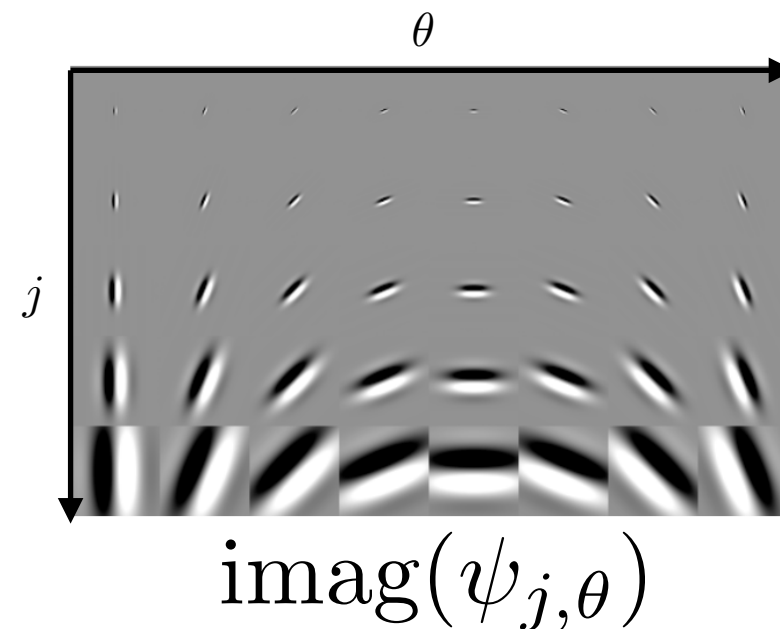
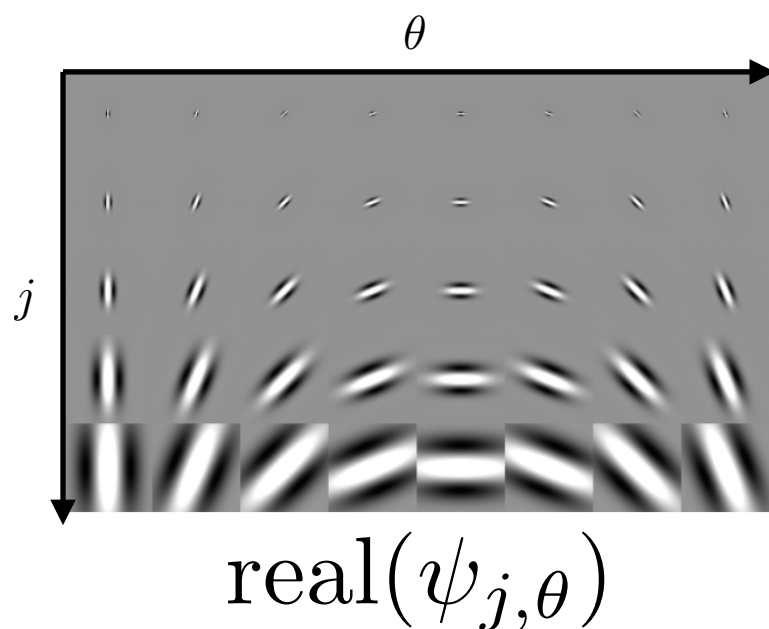


$$|r_k - r_l| \gg 0$$

Wavelets

- Wavelet $\psi(u) = e^{-|u|^2} (e^{i\eta_0 \cdot u} - C)$, $\int \psi = 0$
- Dilated and rotated:

$$\psi_{j,\theta} = 2^{-3j} \psi(2^{-j} R_\theta^{-1} u), \quad j \in \mathbb{Z}, \quad R_\theta \in O(3)$$



Wavelet Transform: $\rho \mapsto \{\int \rho, \quad \rho * \psi_{j,\theta}(u)\}$

Interactions at scale 2^j in direction θ

Wavelet modulus

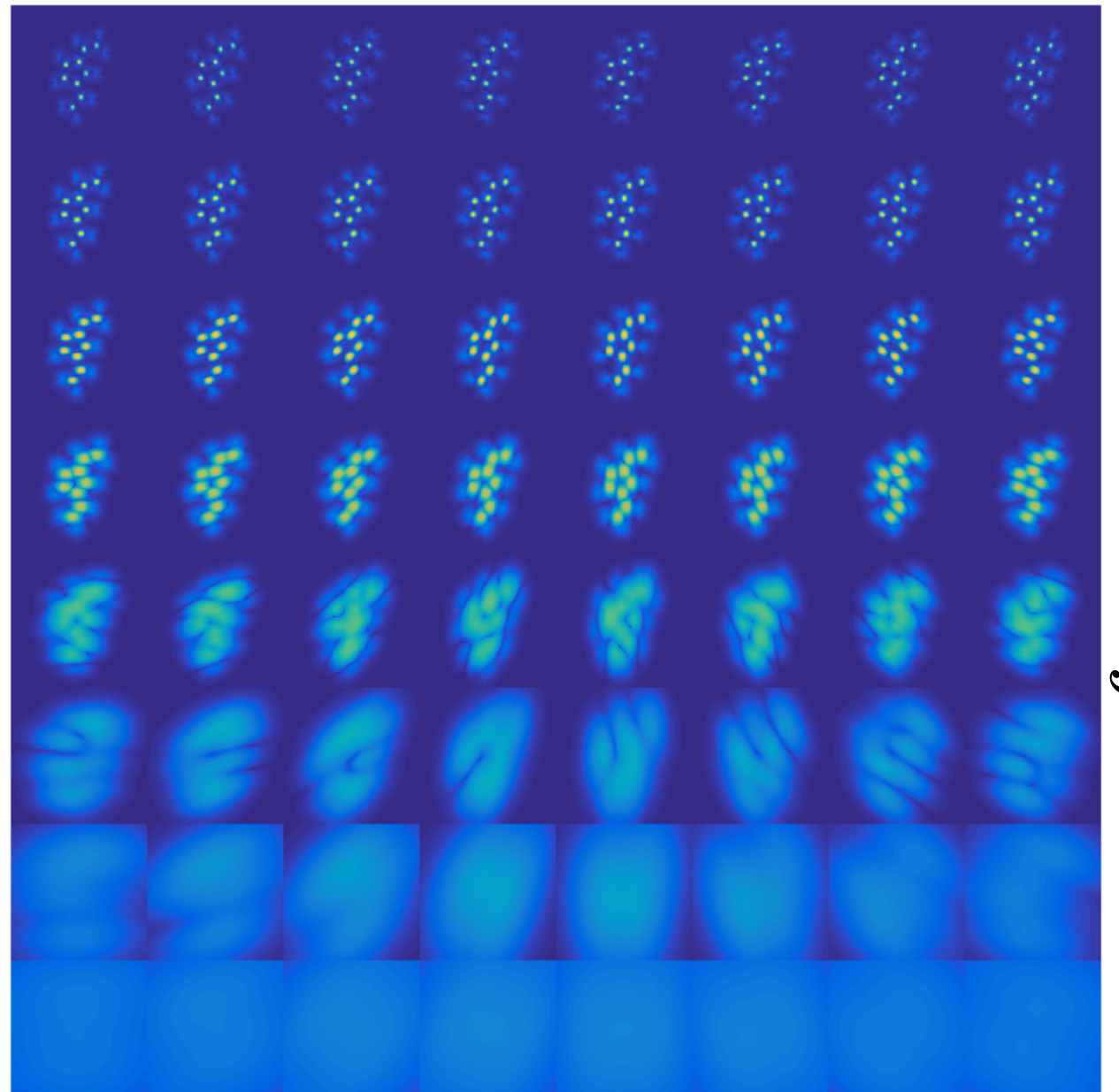
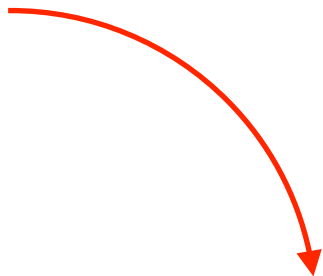
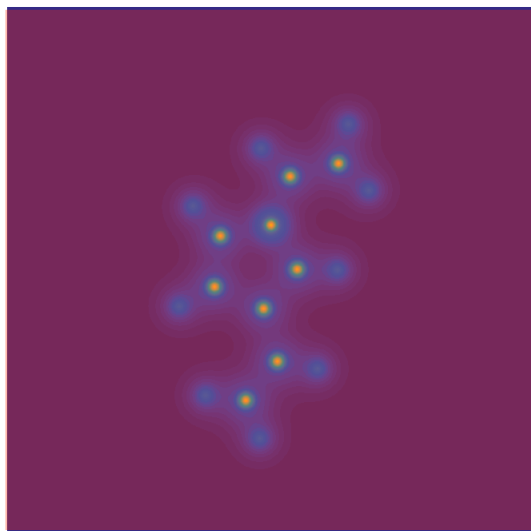
Rotations θ

Scales j

$$|\rho * \psi_{j,\theta}(u)|$$

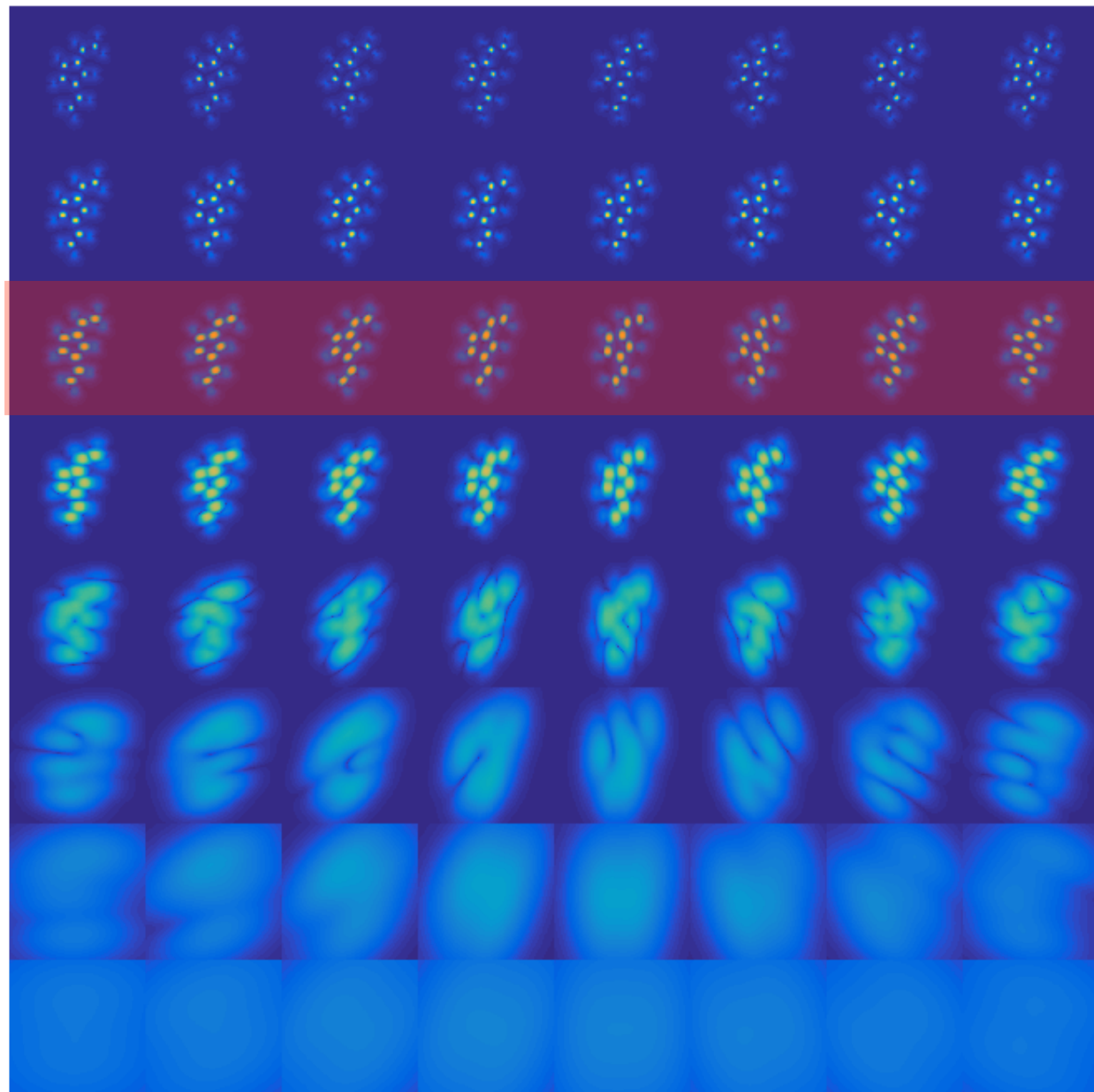
$\rho(u)$

$$\|\rho\|_1 = \hat{\rho}(0)$$



Wavelet invariants

Rotations θ



Scales

$$|\rho * \psi_{j,\theta}(u)|$$

$$\begin{aligned} & \| \rho * \psi_{j,\cdot} \|_1 \\ &= \int_{\mathrm{O}(3)} \int_{\mathbb{R}^3} | \rho * \psi_{j,\theta}(u) | \, du \, d\theta \end{aligned}$$

Invariant wavelet operator

Invariant wavelet operator:

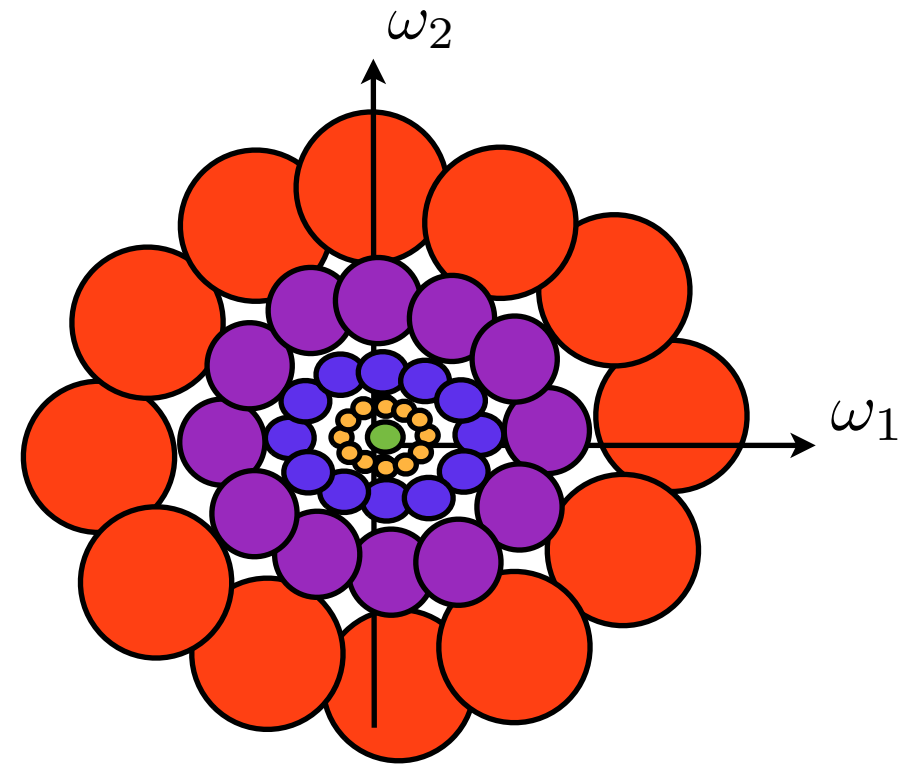
$$\Theta\rho(j) = \|\rho * \psi_{j,\cdot}\|_1$$

Pros:

- Isometry invariant
- Stable to deformations (*Mallat*)
- Diagonalizes Coulomb and is sparse:

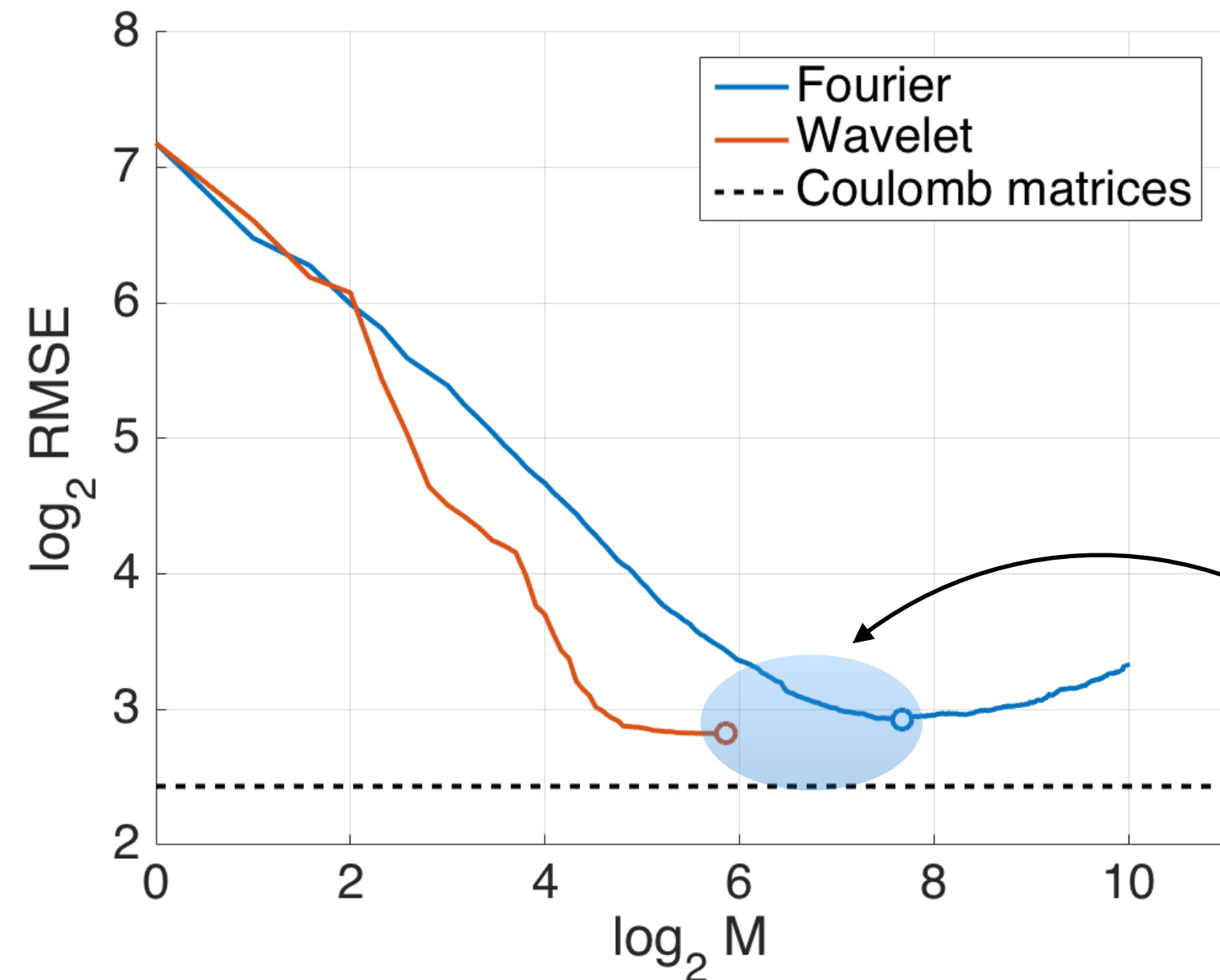
$$U(\rho) = \sum_{j=\log \epsilon}^{-\log \epsilon} w_j \|\rho * \psi_{j,\cdot}\|_2^2 + O(\epsilon)$$

Cons: Encoding the invariants removes a lot of information

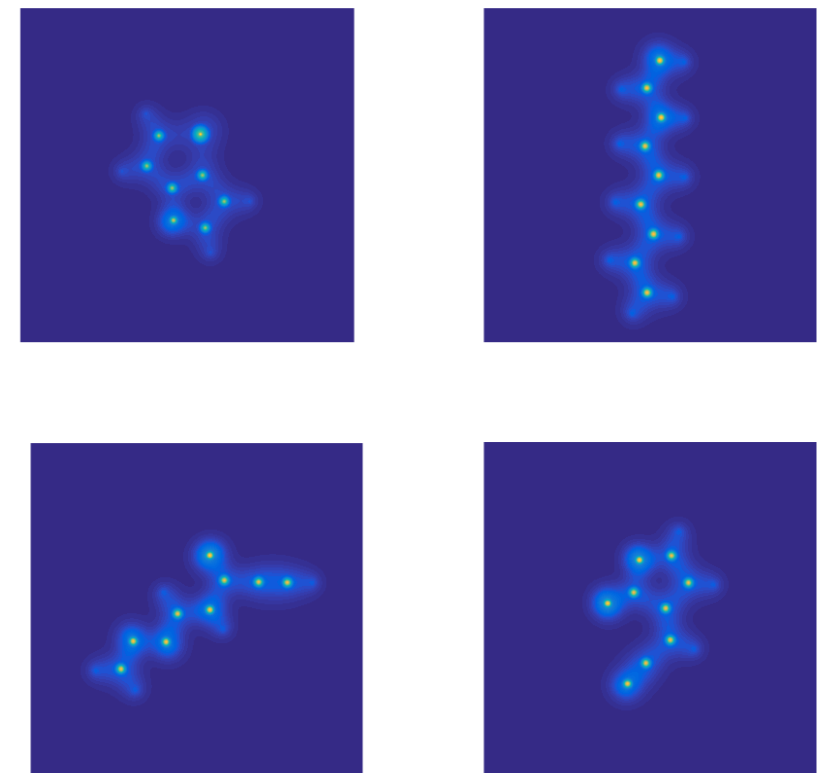


M-term regression error

$$\tilde{f}_M(x) = \sum_{m=1}^M w_{k_m} \phi_{k_m}(x)$$



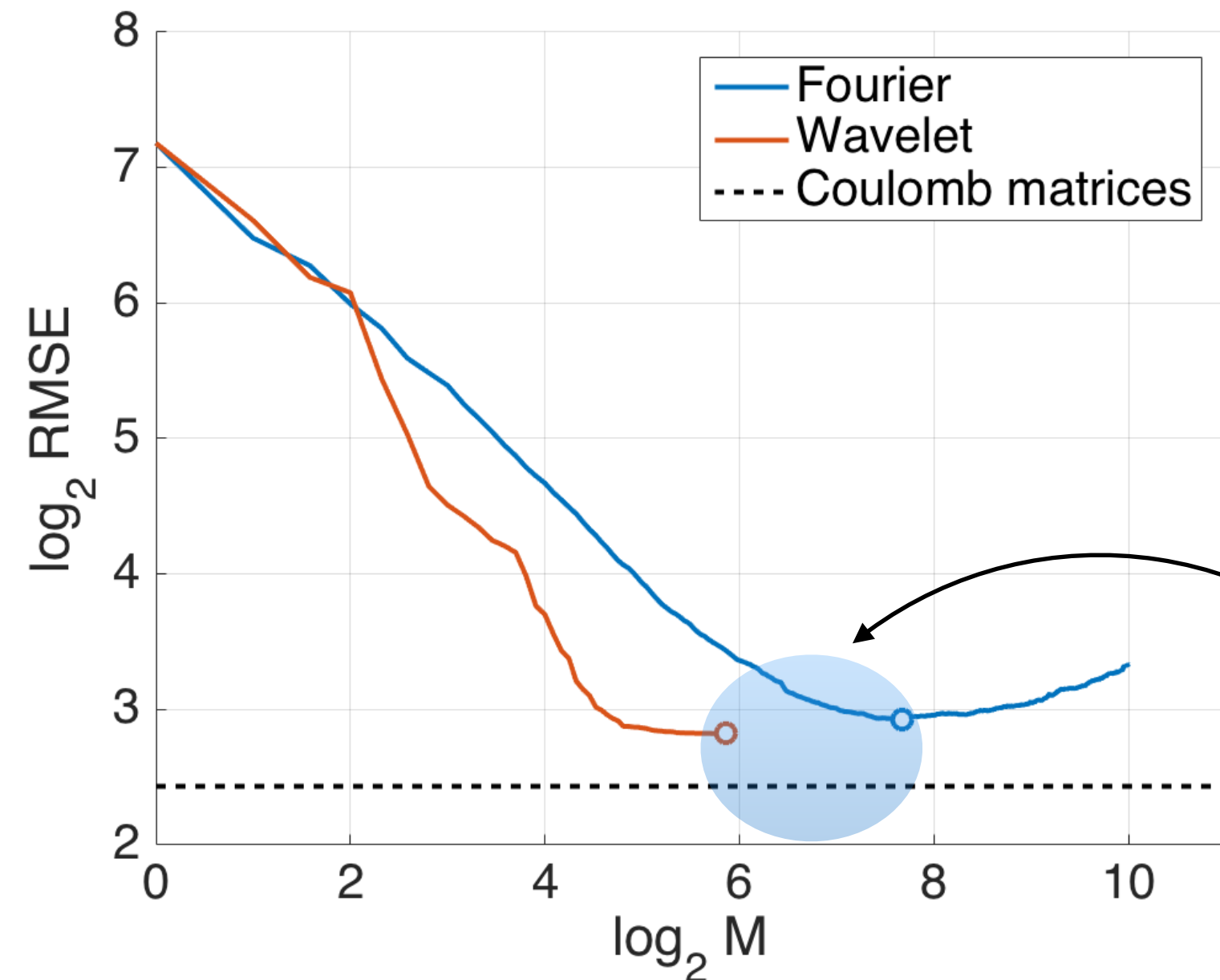
Data set:
Planar organic molecules



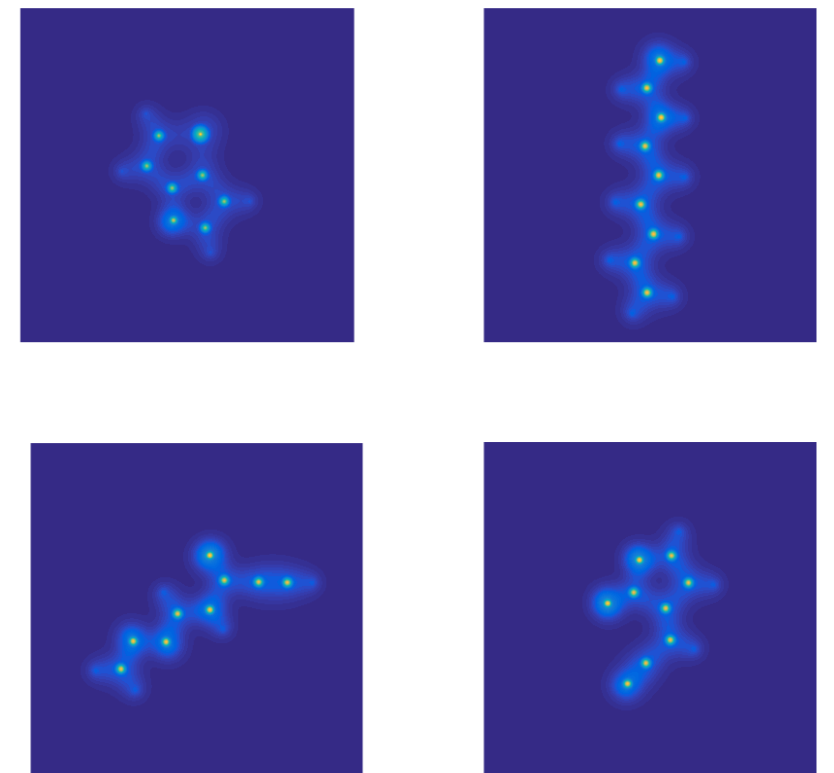
Multiscale wavelet dictionary
is sparser than Fourier dictionary

M-term regression error

$$\tilde{f}_M(x) = \sum_{m=1}^M w_{k_m} \phi_{k_m}(x)$$

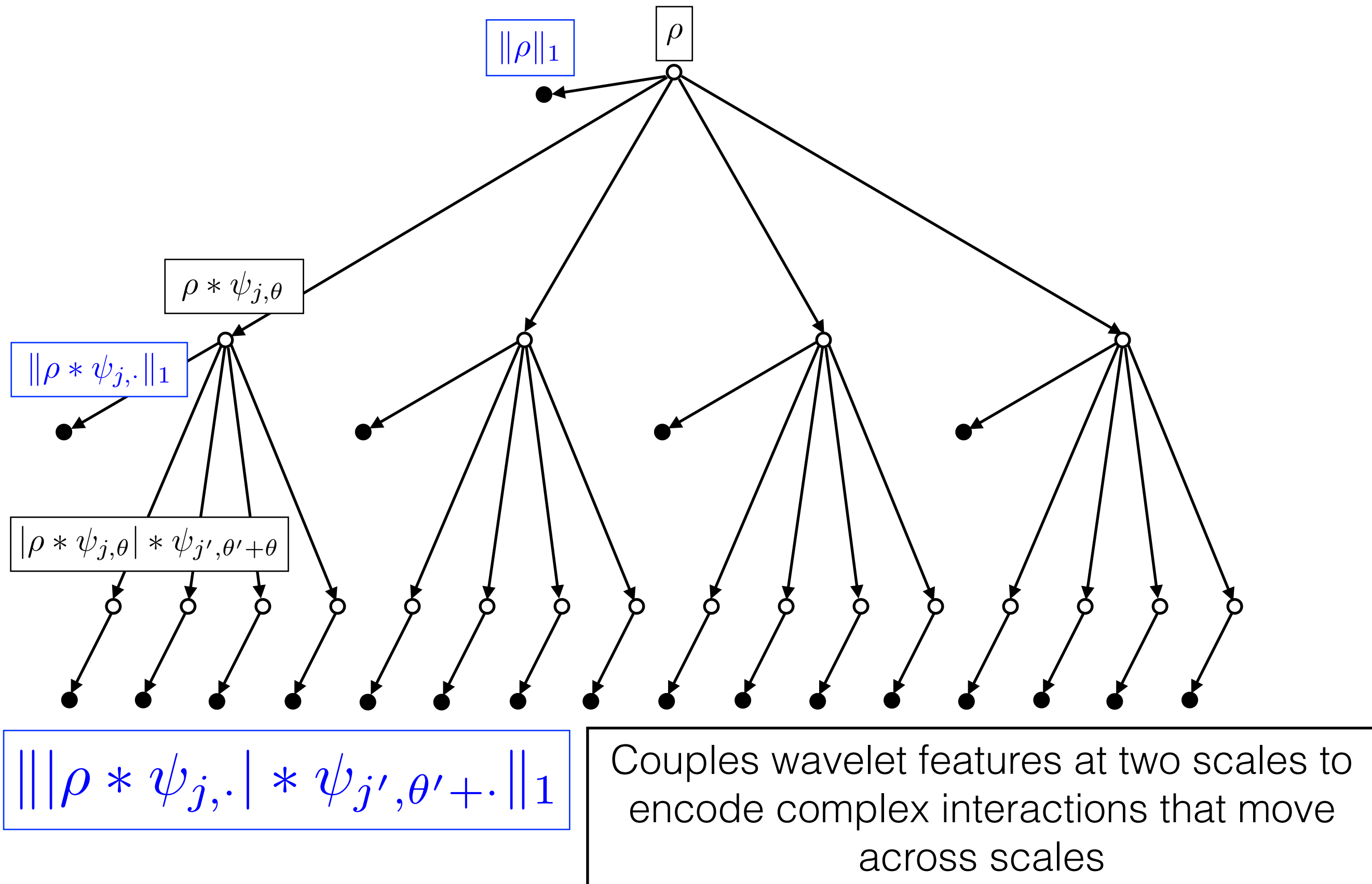


Data set:
Planar organic molecules

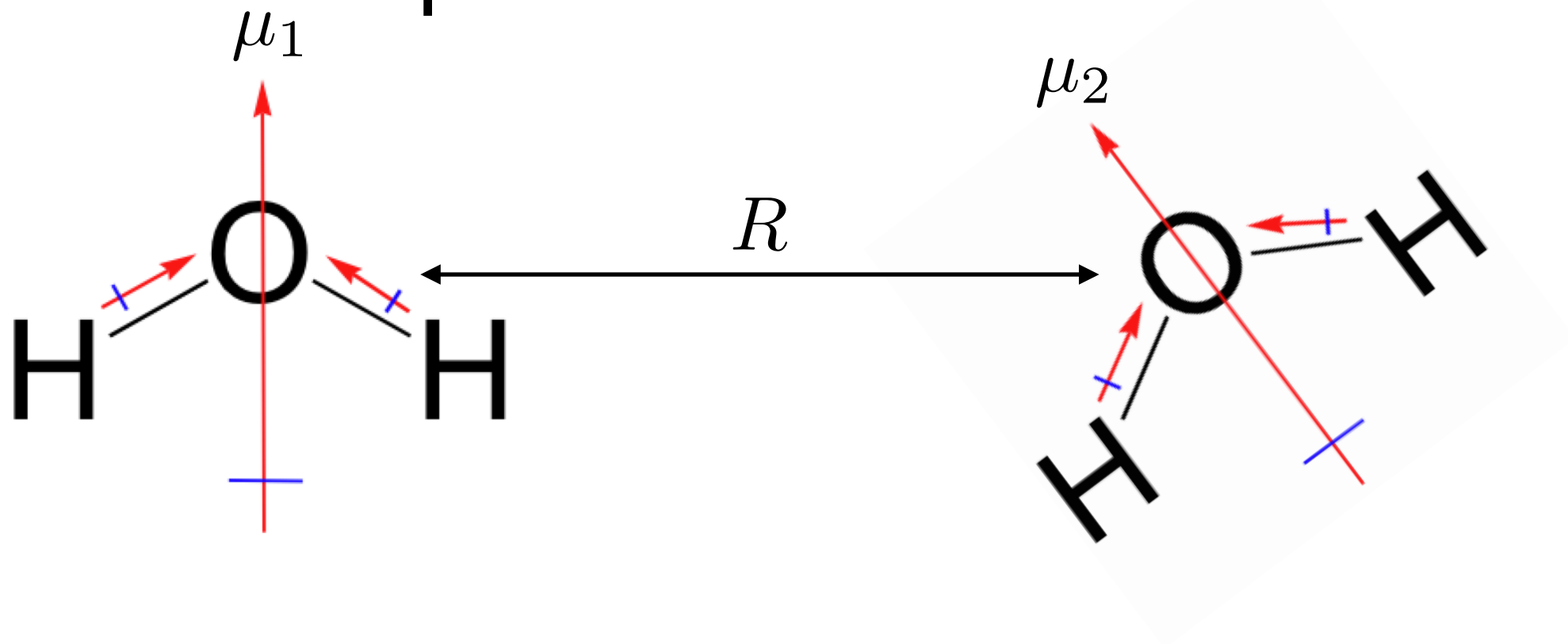


Coulomb matrices outperform
both Fourier and wavelets

Invariant scattering operator



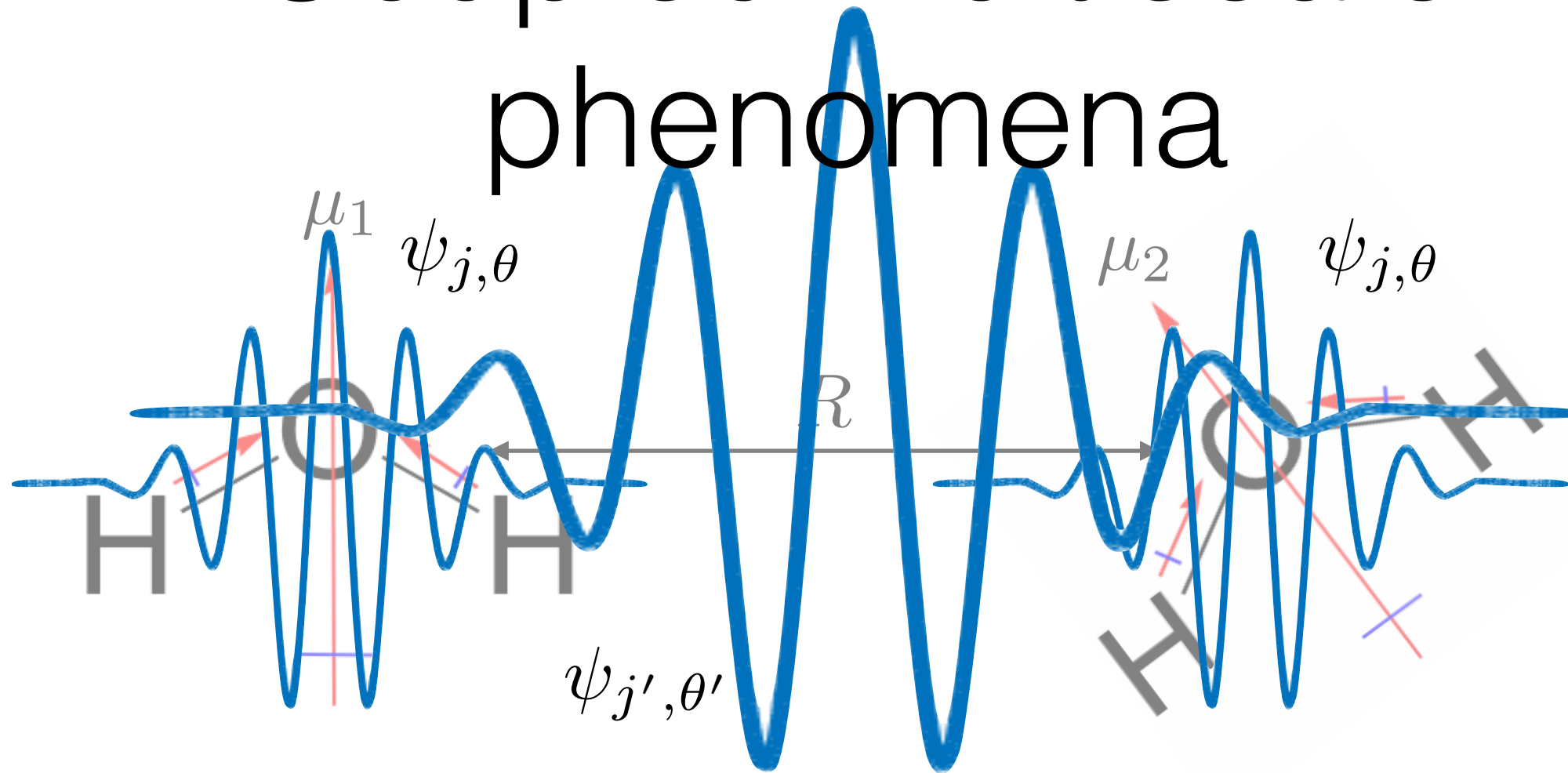
Coupled multiscale phenomena



Multipole expansion dipole-dipole moment:

$$\frac{\mu_1 \cdot \mu_2}{R^3}$$

Coupled multiscale phenomena



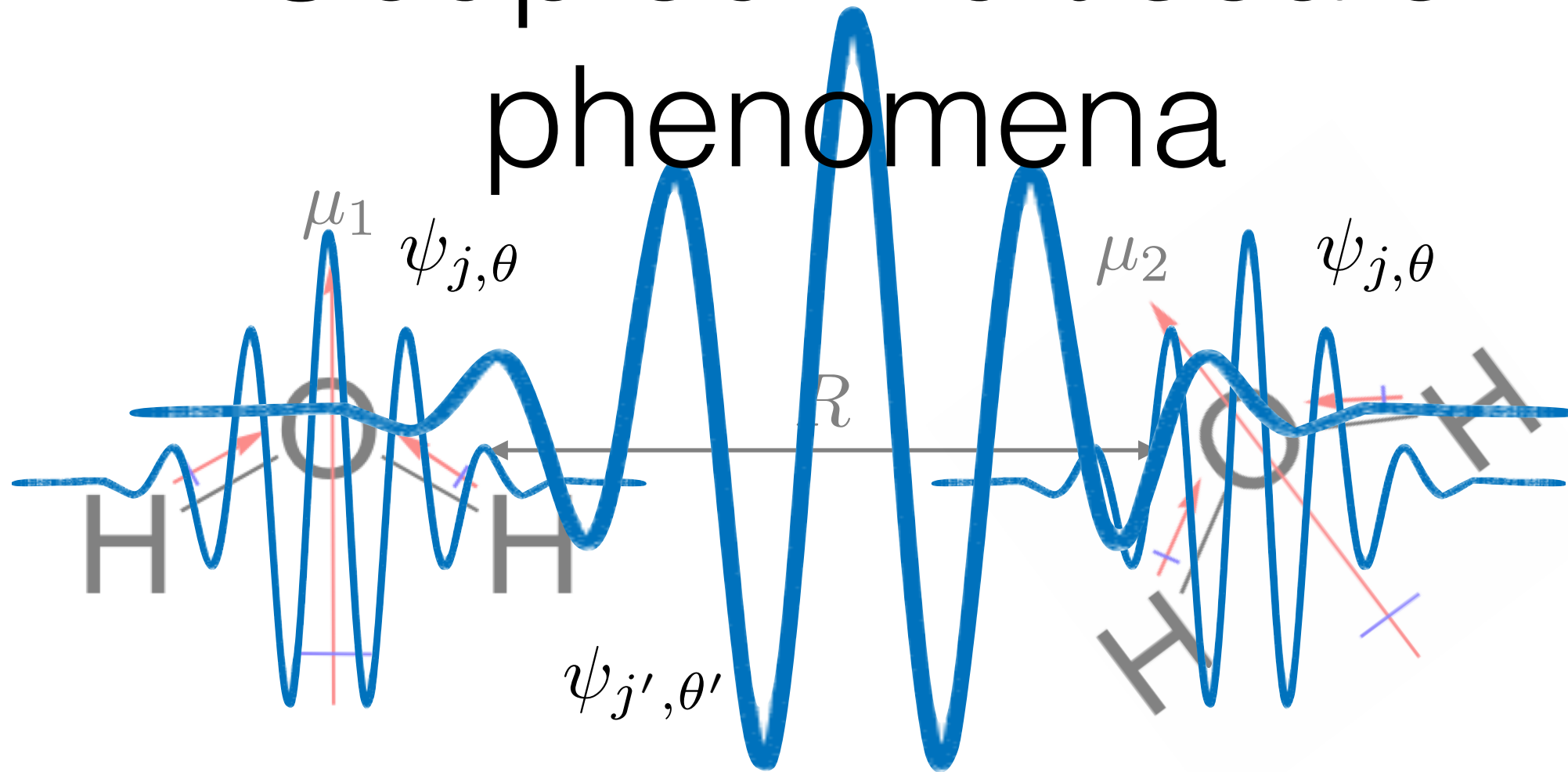
Small scale wavelets $\psi_{j,\theta}$ can learn dipole orientations

Large wavelets $\psi_{j',\theta'}$ can learn the distance between molecules

But there is no linear expansion of the dipole-dipole moment as:

$$\frac{\mu_1 \cdot \mu_2}{R^3} \neq \sum_{j \text{ small}} w_j \|\rho * \psi_{j,\cdot}\|_1 + \sum_{j' \text{ large}} w_{j'} \|\rho * \psi_{j',\cdot}\|_1$$

Coupled multiscale phenomena



2nd order scattering features on the other hand:

First apply the small scale wavelet transform $\rho * \psi_{j,\theta} \dots$

\dots then apply the large scale wavelet transform $|\rho * \psi_{j,\theta}| * \psi_{j',\theta'+\theta}$

Intuition and numerical evidence supports (theory in progress):

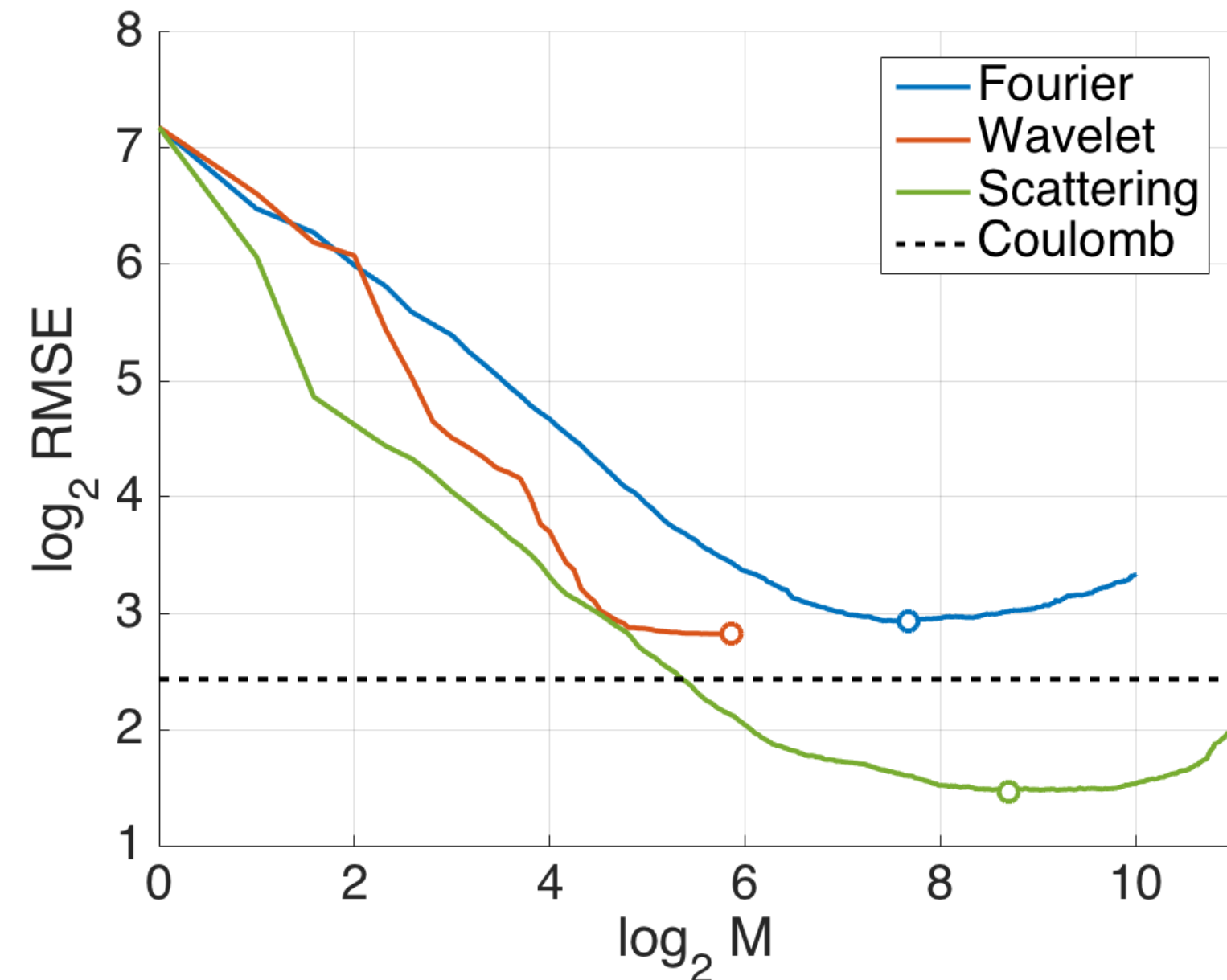
$$\frac{\mu_1 \cdot \mu_2}{R^3} \approx \sum_{j \text{ small}} \sum_{j' \text{ large}} \sum_{\theta'} w_{j,j',\theta'} |||\rho * \psi_{j,\cdot}| * \psi_{j',\theta'+\cdot}||_1$$

M-term regression error

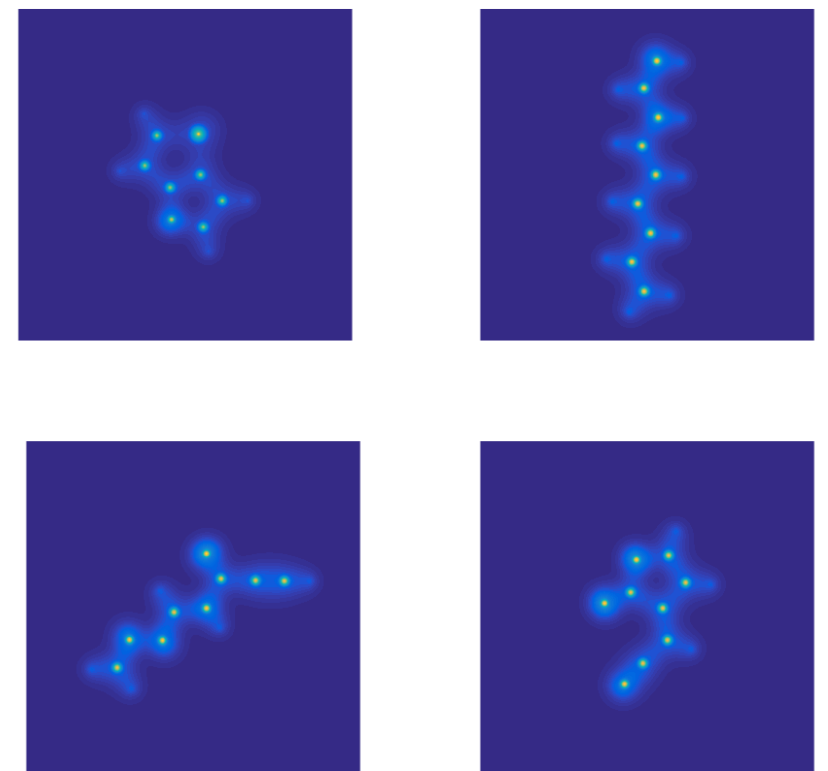
$$\tilde{f}_M(x) = \sum_{m=1}^M w_{k_m} \phi_{k_m}(x)$$

Errors (kcal/mol)

	Coulomb	Fourier	Wavelet	Scattering
MAE	2.4	5.3	5.4	1.7
RMSE	5.4	7.2	7.1	2.6



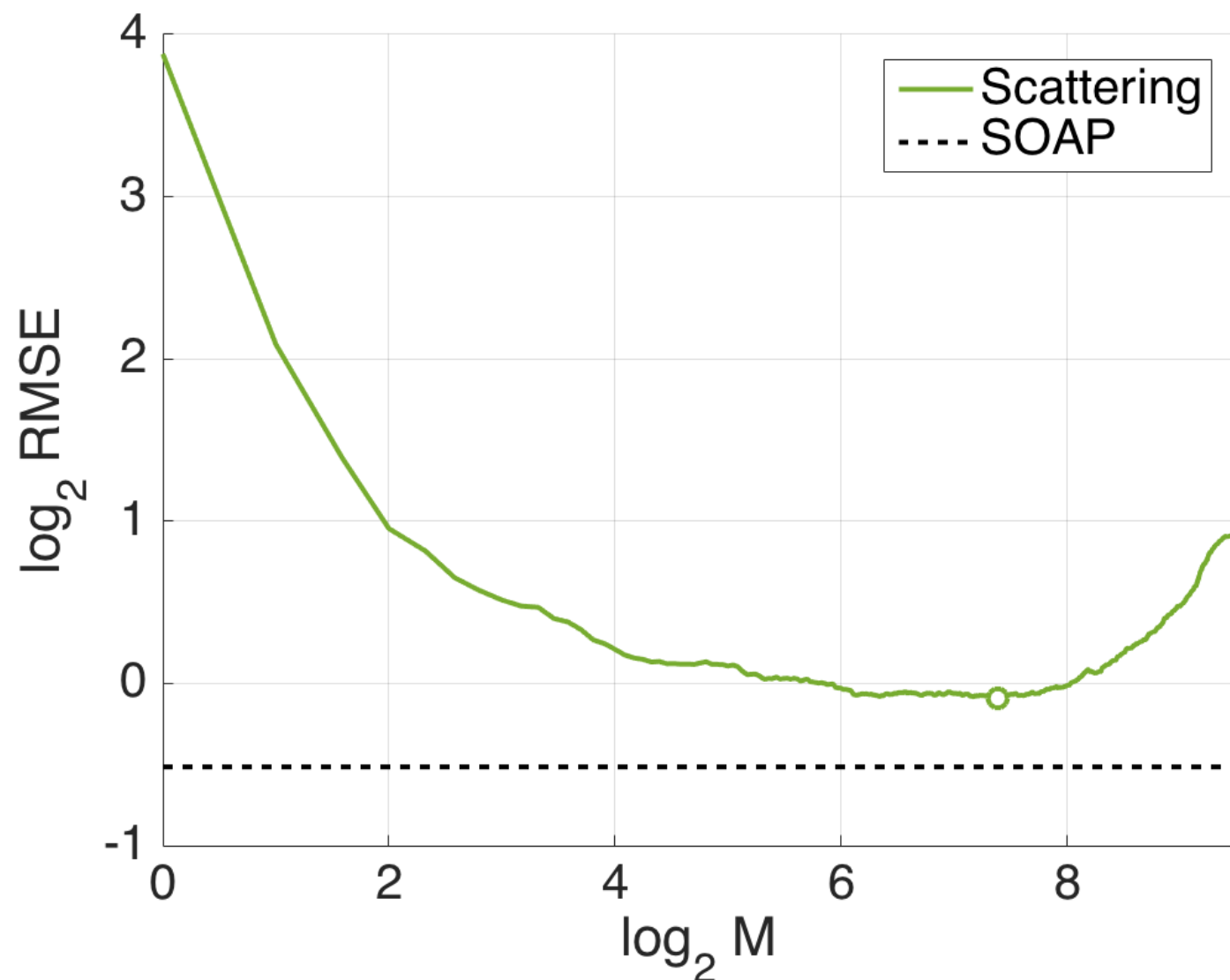
Data set:
Planar organic molecules



Water Molecules

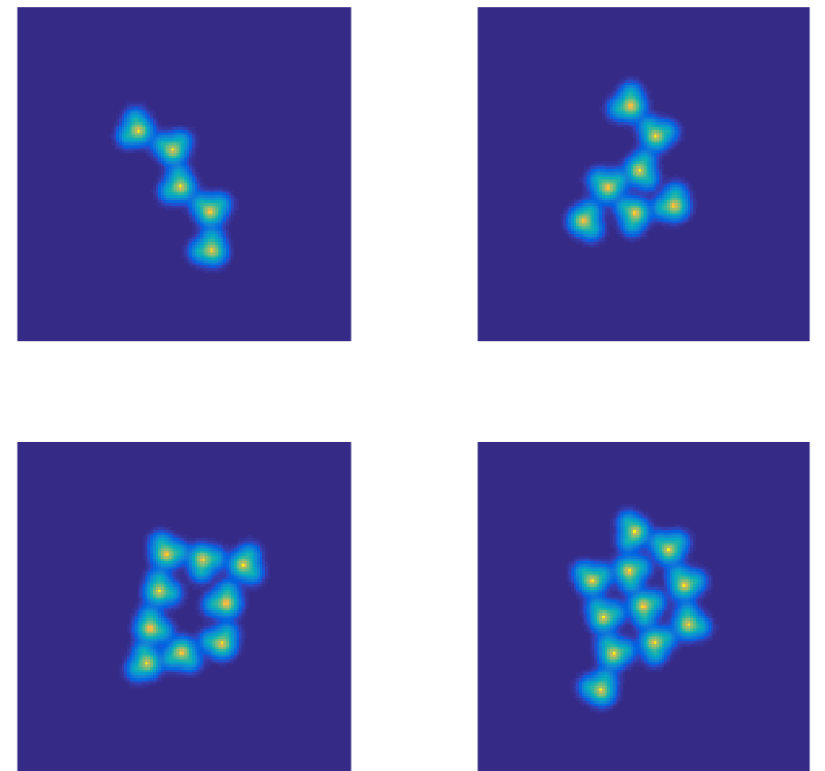
Root mean square error (meV/atom)

	SOAP	Scattering
700 Training	0.70	0.94



Long range interactions are
 $O(1/R^3)$

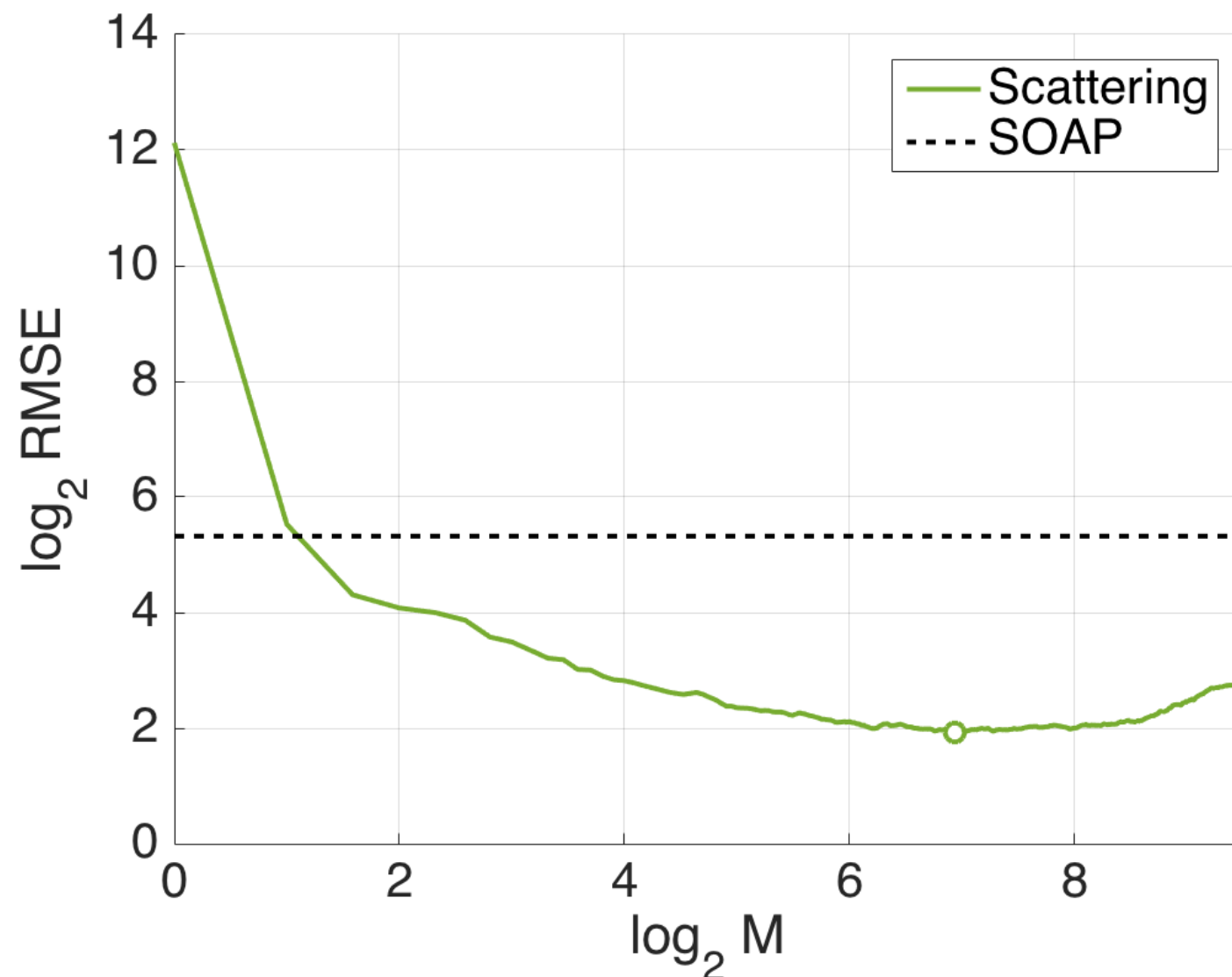
Data set:
Water molecules



Water Lithium Ion Systems

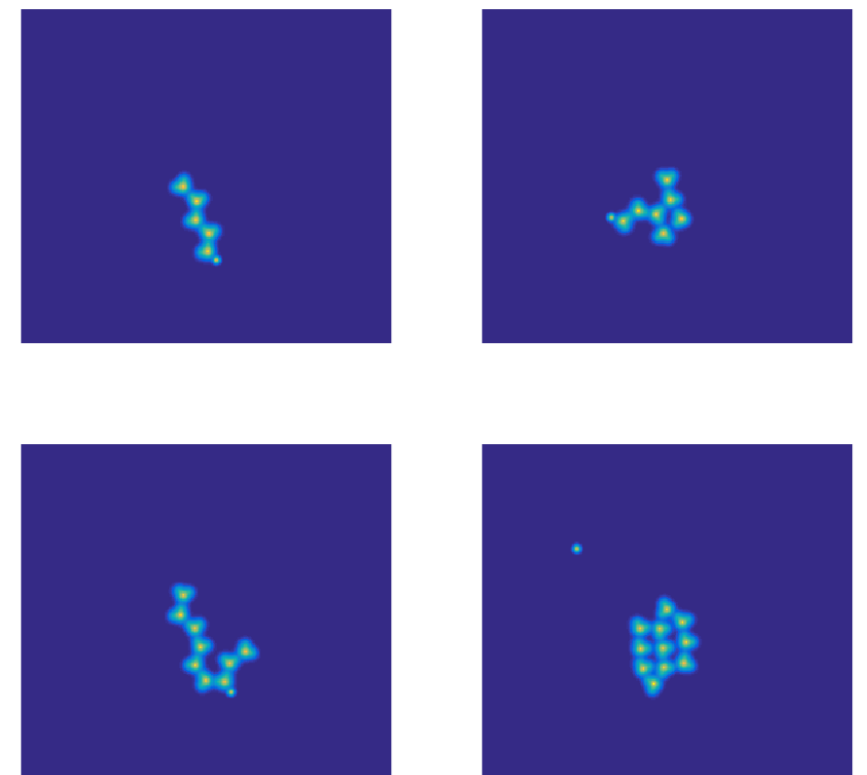
Root mean square error (meV/atom)

	SOAP	Scattering
700 Training	40	3.8



Long range interactions are
 $O(1/R^2)$

Data set:
Water & Lithium Ion



Future directions

- Efficient 3D software
- Rigorous link of scattering second layer with multipole methods
- Theoretical analysis of kinetic and exchange-correlation energies
- Utilization of forces
- Extensions to solid state physics, drug discovery, other many body problems in physics, other physical systems exhibiting complex multiscale behavior
- Connections with deep learning

Conclusions

To apply machine learning to many body problems, we must attack the curse of dimensionality. To do so we need to:

- Interpolate in low dimensional approximation spaces by...
- ...nonlinear separation of original variables into learned non-interacting variables.

To achieve these goals, we utilize:

- Scale and angular separation over groups (translation, rotation)
- Symmetry and invariance properties
- Stability to deformations
- Cascade of operators (wavelet modulus) with the previous three properties to recover lost information